

# Advance Topics in Astrophysical Fluids

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*Numerical Astrophysics Summer School 2019:  
Astrophysical Fluid Dynamics*

# Multi-dimension

- **Directional splitting:**  $\frac{\partial U}{\partial t} + \frac{\partial F_x}{\partial x} = 0 \rightarrow \frac{\partial U}{\partial t} + \frac{\partial F_y}{\partial y} = 0 \rightarrow \frac{\partial U}{\partial t} + \frac{\partial F_z}{\partial z} = 0$ 
  - **Swap directions in the next step to improve accuracy**
    - E.g.,  $(x, y, z) \rightarrow (z, y, x)$
  - **Pro: more stable (in general), larger timestep**  $\Delta t \leq \frac{\Delta h}{\max(|v_x| + |v_y| + |v_z|) + C_s}$
  - **Con: break spatial symmetry**
- **Directional unsplitting:**  $\frac{\partial U}{\partial t} + \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z} = 0$ 
  - **Both Riemann solver and data reconstruction are still 1D**
  - **Pro: preserve spatial symmetry**
  - **Con: may be less stable (i.e., negative density/pressure), smaller timestep**
    - E.g., MUSCL-Hancock:  $\Delta t \leq \frac{\Delta h}{|v_x| + |v_y| + |v_z| + 3C_s}$

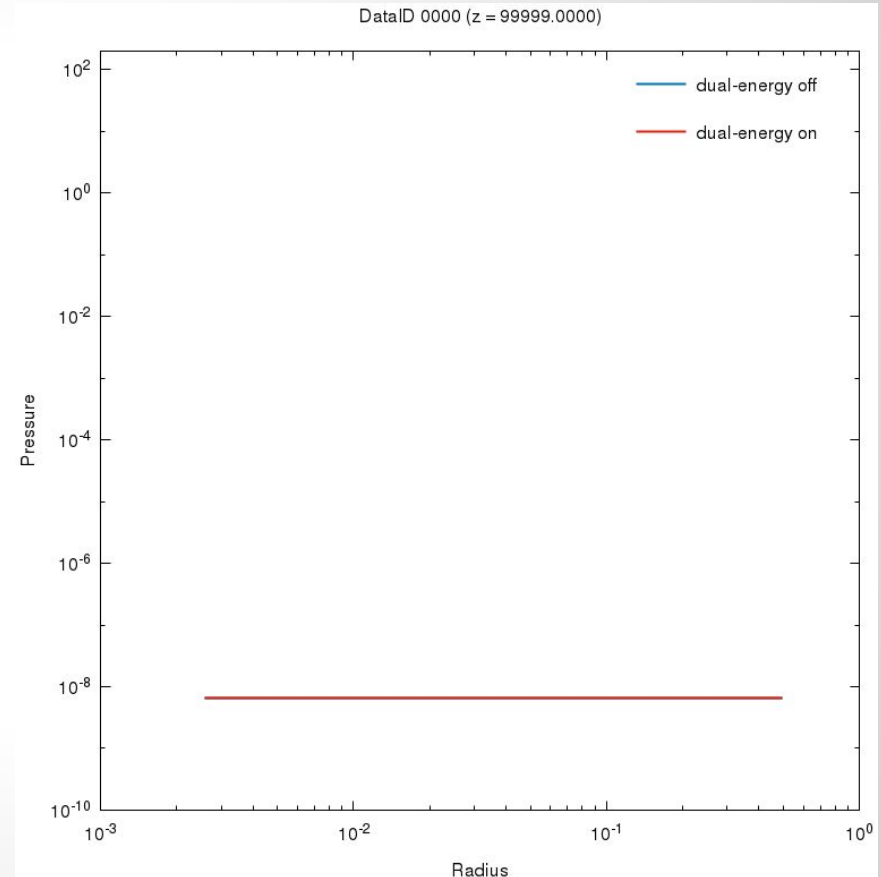
# Issue of Negative Pressure

- **Ideal gas:**  $P = (\gamma - 1)e = (\gamma - 1)(E - E_k) = (\gamma - 1) \left[ E - \frac{(\rho v)^2}{2\rho} \right]$ 
  - **Conserved variables**  $E, \rho, \rho v$  **suffer from truncation errors**
  - **When**  $E \approx E_k \gg e$  **(i.e., high-speed cold flow),**  $e$  **may be seriously contaminated by truncation errors and become negative**
  - **Irrelevant for dynamics, but crucial if temperature is required**
    - **E.g., cooling rate may depend on temperature**
- **Popular solution: dual-energy formalism**
  - **Solve an additional auxiliary eq. → either entropy**  $s$  **or internal energy**  $e$ 

$$\frac{\partial s}{\partial t} + \nabla \cdot (s\mathbf{v}) = 0 \quad \text{or} \quad \frac{\partial e}{\partial t} + \nabla \cdot (e\mathbf{v}) = -P\nabla \cdot \mathbf{v}$$
  - **Use**  $s$  **or**  $e$  **to compute**  $P$  **only in the regions without shocks**
    - **This is the most tricky part**

# Issue of Negative Pressure

- **Example: spherical collapse in cosmology**
- **Solving entropy in this example**
- **Dual-energy formalism significantly improves the solution in the pre-shock (upstream) region**
- **Post-shock (downstream) region is not affected → a strong shock is still captured**



# Self-gravity

$$\begin{aligned}\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0 \\ \frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} + P \mathbf{I}) &= \underline{-\rho \nabla \phi} \\ \frac{\partial E}{\partial t} + \nabla \cdot [(E + P) \mathbf{v}] &= \underline{-\rho \mathbf{v} \cdot \nabla \phi}\end{aligned}$$



Poisson eq.

$$\nabla^2 \phi = 4\pi G \rho$$

- **Operator *splitting* method:**
  - **Procedure**
    - Solve the original Euler eqs. without gravity
    - Solve the Poisson equation for  $\Phi$
    - Update  $\rho \mathbf{v}$  and  $E$  by gravity without considering hydro fluxes
  - **Disadvantage**
    - Inadequate for handling hydrostatic equilibrium
    - Only 1<sup>st</sup>-order accurate

# Self-gravity

- Operator *unsplitting* method:
  - Procedure
    - Solve the Poisson equation for  $\phi^n$
    - Solve the Euler eqs. with the half-step velocity corrected by gravity
$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla \phi$$
    - Solve the Poisson equation again for  $\phi^{n+1}$
    - Correct  $\rho v$  and  $E$  with the half-step gravity  $\phi^{n+1/2} \approx (\phi^n + \phi^{n+1})/2$
  - Advantage
    - Better for handling hydrostatic equilibrium
    - 2<sup>nd</sup>-order accurate
- Both methods do NOT conserve total momentum and energy in general
  - Cause: gravity is treated as source terms
  - Solution: rewrite  $\rho \nabla \phi$  and  $\rho \mathbf{v} \cdot \nabla \phi$  into flux-conservative forms

# Magnetohydrodynamics (MHD)

- Ideal MHD:
  - $$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad \leftarrow \text{mass conservation}$$
  - $$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + P^* \mathbf{I}) = 0 \quad \leftarrow \text{momentum conservation}$$
  - $$\frac{\partial E}{\partial t} + \nabla \cdot [(E + P^*) \mathbf{v} - \mathbf{B}(\mathbf{B} \cdot \mathbf{v})] = 0 \quad \leftarrow \text{energy conservation}$$
  - $$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{v} \times \mathbf{B}) = 0 \quad \leftarrow \text{induction eq. + ideal Ohm's law}$$
- $E = e + \frac{1}{2} \rho v^2 + \frac{B^2}{2}, \quad P^* = P + \frac{B^2}{2}$
- 9 variables to be solved by the 8 equations above + equation of state
- Divergence-free constraint on the magnetic field:  $\nabla \cdot \mathbf{B} = 0$

# Flux-conservative Form for MHD

- $$\frac{\partial U}{\partial t} + \frac{\partial \mathbf{F}_x}{\partial x} + \frac{\partial \mathbf{F}_y}{\partial y} + \frac{\partial \mathbf{F}_z}{\partial z} = 0,$$

$$U = \begin{bmatrix} \rho \\ \rho v_x \\ \rho v_y \\ \rho v_z \\ E \\ B_x \\ B_y \\ B_z \end{bmatrix}, \quad \mathbf{F}_x = \begin{bmatrix} \rho v_x \\ \rho v_x^2 + P^* - B_x^2 \\ \rho v_x v_y - B_x B_y \\ \rho v_x v_z - B_x B_z \\ (E + P^*)v_x - B_x (\mathbf{B} \cdot \mathbf{v}) \\ 0 \\ v_x B_y - v_y B_x \\ v_x B_z - v_z B_x \end{bmatrix}, \text{ similarly for } \mathbf{F}_y, \mathbf{F}_z$$

- Fluid conserved variables can be updated similarly using the finite-volume scheme for pure hydro**
- Key question: how to update the magnetic field and ensure the divergence-free constraint?**



# Constrained Transport (CT) Method

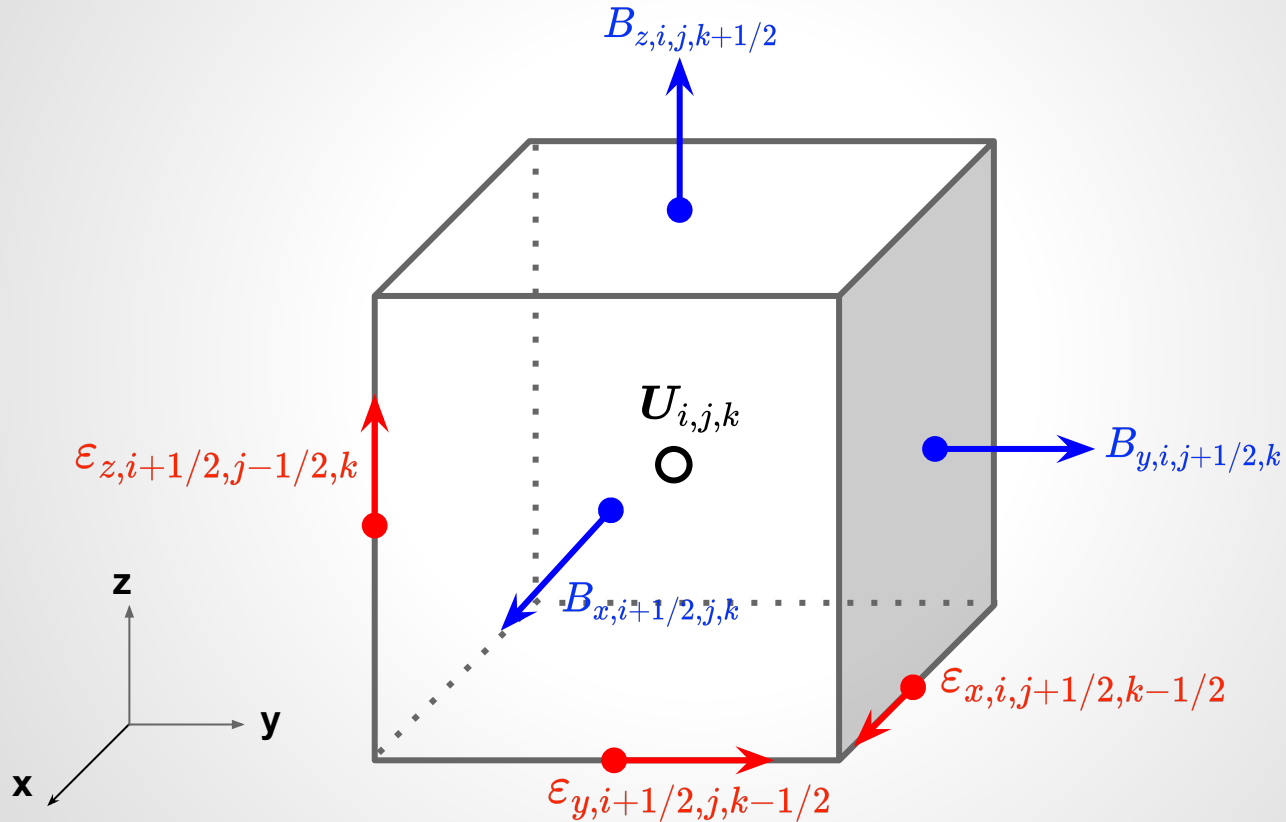
- **Stokes' theorem:** 
$$\int_A \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{A} = \int_A [\nabla \times (\mathbf{v} \times \mathbf{B})] \cdot d\mathbf{A} = \oint_{\partial A} \mathbf{v} \times \mathbf{B} \cdot d\mathbf{l}$$
  - **Electromotive force (EMF):**  $\boldsymbol{\varepsilon} = -\mathbf{v} \times \mathbf{B}$
- **Integrate over cell area (e.g.,  $\Delta y \Delta z$ ) and time interval  $\Delta t = t^{n+1} - t^n$**

$$B_{x,i-1/2,j,k}^n \equiv \frac{1}{\Delta y \Delta z} \int_{z_{k-1/2}}^{z_{k+1/2}} \int_{y_{j-1/2}}^{y_{j+1/2}} B_x(x_{i-1/2}, y, z, t^n) dy dz$$
$$\varepsilon_{y,i-1/2,j,k-1/2}^{n+1/2} \equiv \frac{1}{\Delta y \Delta t} \int_{t^n}^{t^{n+1}} \int_{y_{j-1/2}}^{y_{j+1/2}} \varepsilon_y(x_{i-1/2}, y, z_{k-1/2}, t) dy dt$$
$$\varepsilon_{z,i-1/2,j-1/2,k}^{n+1/2} \equiv \frac{1}{\Delta z \Delta t} \int_{t^n}^{t^{n+1}} \int_{z_{k-1/2}}^{z_{k+1/2}} \varepsilon_z(x_{i-1/2}, y_{j-1/2}, z, t) dz dt$$

# Constrained Transport (CT) Method

- $$B_{x,i-1/2,j,k}^{n+1} = B_{x,i-1/2,j,k}^n - \frac{\Delta t}{\Delta y} \left( \varepsilon_{z,i-1/2,j+1/2,k}^{n+1/2} - \varepsilon_{z,i-1/2,j-1/2,k}^{n+1/2} \right) + \frac{\Delta t}{\Delta z} \left( \varepsilon_{y,i-1/2,j,k+1/2}^{n+1/2} - \varepsilon_{y,i-1/2,j,k-1/2}^{n+1/2} \right)$$
- This form is again exact → similar to the finite-volume formulation
- $B_{x,i-1/2,j,k}^n$  : area-averaged magnetic field
- $\varepsilon_{z,i-1/2,j\pm 1/2,k}^{n+1/2}$ ,  $\varepsilon_{y,i-1/2,j,k\pm 1/2}^{n+1/2}$  : time- and line-averaged EMF
- Similar expressions can be derived for  $B_{y,i,j-1/2,k}^{n+1}$  &  $B_{z,i,j,k-1/2}^{n+1}$
- Area-averaged magnetic field are located at the cell faces instead of centers → staggered grid

# Staggered Grid in CT



# Divergence Free in CT

- Finite-volume representation of the divergence-free constraint:

$$\frac{1}{\Delta x \Delta y \Delta z} \int_{V_{i,j,k}} (\nabla \cdot B^n) dV = 0$$

$$\rightarrow (\nabla \cdot B^n)_{i,j,k} = \frac{B_{x,i+1/2,j,k}^n - B_{x,i-1/2,j,k}^n}{\Delta x} + \frac{B_{y,i,j+1/2,k}^n - B_{y,i,j-1/2,k}^n}{\Delta y} + \frac{B_{z,i,j,k+1/2}^n - B_{z,i,j,k-1/2}^n}{\Delta z} = 0$$

exact form

HW

- CT update guarantees  $\nabla \cdot B^{n+1} = \nabla \cdot B^n$ 
  - **Divergence-free constraint is preserved to the machine precision assuming that it is satisfied at the beginning**
  - The exact way to compute EMF varies from scheme to scheme

# Multi-species

- In astrophysical environments, it is common that there are more than one species. From species  $i$  to  $N$ , each species follows,

$$\frac{\partial \rho_i}{\partial t} + \nabla \cdot (\rho_i v) = m_i \Omega_i$$

$\Omega_i$  Number density production rate for specie  $i$

$$\sum_i m_i \Omega_i = 0$$

- Summing the species equations gives the continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho v) = 0$$

# **Advanced usage of FLASH code**

# Advanced usage of FLASH code

- **Dimension and Geometry**
- **Split and Unsplit solvers**
- **Dual energy**
- **Self-gravity**
- **MHD**
- **Multi-species**
- **EoS**
- **Modify your simulation**
- **Modify AMR rules**

# Advanced usage of FLASH code

- Dimension and Geometry

```
./setup -1d +spherical
```

```
./setup -2d +cylindrical
```

```
./setup -3d +pm4dev
```

```
./setup -3d +ug
```

**Table 8.7:** Different geometry types. For each geometry/dimensionality combination, the "support" column lists the `GRID` implementations that support it: `pm4` stands for `PARAMESH 4.0` and `PARAMESH 4dev`, `pm2` for `PARAMESH 2`, `UG` for Uniform Grid implementations, respectively.

name	dimensions	support	axisymmetric	X -coord	Y -coord	Z -coord
cartesian	1	pm4,pm2,UG	n	$x$		
cartesian	2	pm4,pm2,UG	n	$x$	$y$	
cartesian	3	pm4,pm2,UG	n	$x$	$y$	$z$
cylindrical	1	pm4,UG	y	$r$		
cylindrical	2	pm4,pm2,UG	y	$r$	$z$	
cylindrical	3	pm4,UG	n	$r$	$z$	$\phi$
spherical	1	pm4,pm2,UG	y	$r$		
spherical	2	pm4,pm2,UG	y	$r$	$\theta$	
spherical	3	pm4,pm2,UG	n	$r$	$\theta$	$\phi$
polar	1	pm4,UG	y	$r$		
polar	2	pm4,pm2,UG	n	$r$	$\phi$	
"polar + z " (cylindrical with a different ordering of coordinates)	3	--	n	$r$	$\phi$	$z$



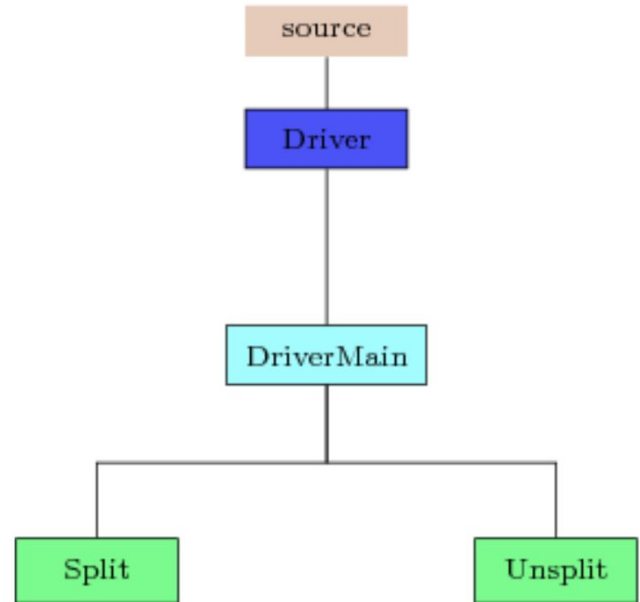
# Split and unsplit solvers

unsplit (default)

```
./setup +uhd
```

Split solver

```
./setup +splitHydro
```



**Figure 7.1:** The Driver unit directory tree.

# Runtime parameters for the unsplit Solver

Variable	Type	Default	Description
order	integer	2	Order of method in data reconstruction: 1st order Godunov (FOG), 2nd order MUSCL-Hancock (MH), 3rd order PPM, 5th order WENO.
RiemannSolver	string	"Roe"	Different choices for Riemann solver. ``LLF (local Lax-Friedrichs)", ``HLL", ``HLLC", ``HYBRID", ``ROE", and ``Marquina"
use_auxEintEqn	bool	True	Turn on/off solving the auxiliary internal energy equation
eintSwitch	real	0	If $\epsilon < \text{eintSwitch} \cdot \frac{1}{2} \mathbf{v} ^2$ , use the internal energy equation to update the pressure
slopeLimiter	string	"vanLeer"	Slope limiter (mc, vanLeer, vanLeer1.5, minmod, hybrid, limited)

# Gravity Units

In setup script,

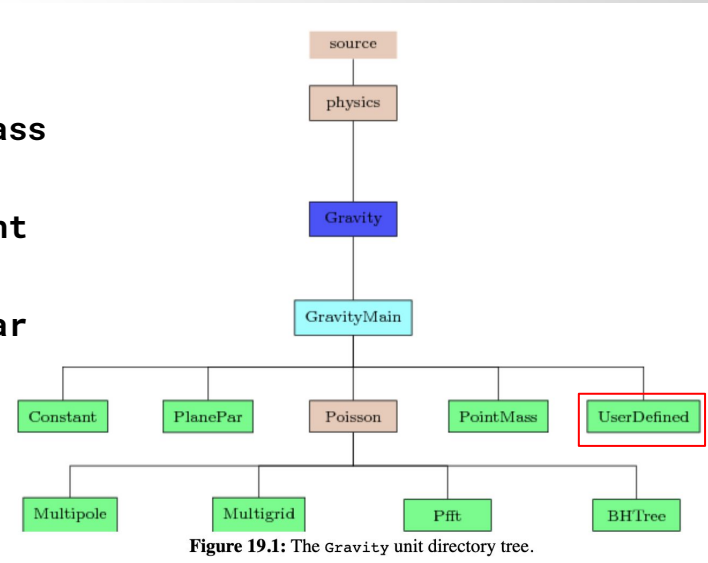
```
./setup ...-unit=source/physics/Gravity/GravityMain/PointMass
```

```
./setup ...-unit=source/physics/Gravity/GravityMain/Constant
```

```
./setup ...-unit=source/physics/Gravity/GravityMain/PlanePar
```

In Config file,

```
REQUIRES physics/Gravity/GravityMain/PointMass
```



# Runtime parameters for gravity

Variable	Type	Default	Description
useGravity	Bool	False	Whether gravity calculations should be performed.
grav_boundary_type	String	"isolated"	Type of gravitational boundary condition if a Poisson solve is used for Gravity; string-valued version of grav_boundary. Accepts: "isolated", "periodic", "dirichlet", and maybe others, depending on the Poisson solver used.

# Runtime parameters for gravity

```
physics/Gravity/GravityMain/Constant
```

```
gconst [REAL] [-981.]
```

```
Valid Values: Unconstrained
```

```
Gravitational acceleration constant
```

```
gdirac [STRING] ["x"]
```

```
Valid Values: Unconstrained
```

```
Direction of acceleration ("x", "y", "z")
```

# Runtime parameters for gravity

physics/Gravity/GravityMain/PlanePar

gravsoft [REAL] [.0001]

Valid Values: Unconstrained  
softening length

ptdirn [INTEGER] [1]

Valid Values: Unconstrained  
x = 1, y = 2, z = 3

ptmass [REAL] [10000.]

Valid Values: Unconstrained  
mass of the point

ptxpos [REAL] [1.]

Valid Values: Unconstrained  
location of the point mass, in the ptdirn direction

# Runtime parameters for gravity

```
physics/Gravity/GravityMain/PointMass
```

```
gravsoft [REAL] [0.001]
```

```
Valid Values: Unconstrained
```

```
ptmass [REAL] [10000.]
```

```
Valid Values: Unconstrained
```

```
ptxpos [REAL] [1.]
```

```
Valid Values: Unconstrained
```

```
ptypos [REAL] [-10.]
```

```
Valid Values: Unconstrained
```

```
ptzpos [REAL] [0.]
```

```
Valid Values: Unconstrained
```

# Self-gravity

`./setup +mpole`

`./setup +gravmpole`

`./setup +gravmgrid`

REQUIRES physics/Gravity/GravityMain/Poisson/Multigrid

REQUIRES physics/Gravity/GravityMain/Poisson/Multipole

## 8.10 GridSolvers

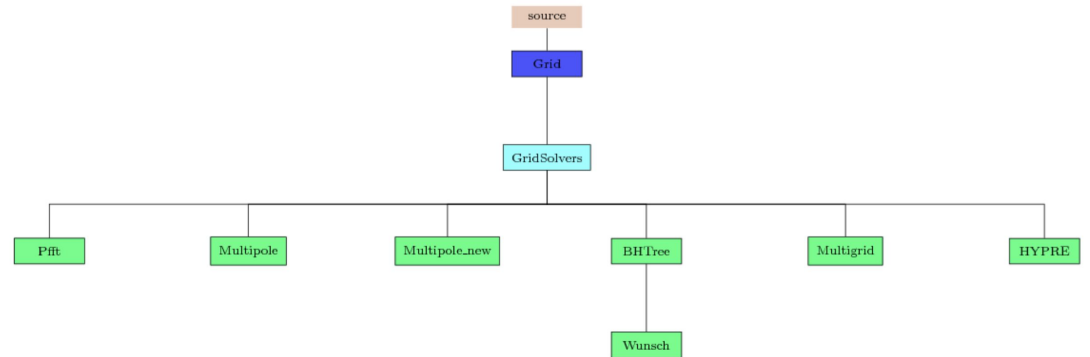


Figure 8.11: The Grid unit: structure of GridSolvers subunit.



# Runtime parameters for gravity

```
physics/Gravity/GravityMain/Poisson
  grav_temporal_extrp [BOOLEAN] [FALSE]
    extrapolate or otherwise rescale
  grav_unjunkPden [BOOLEAN] [TRUE]
    controls whether Gravity_potentialListOfBlocks attempts to restore the
    part of the "pden" ("particle density") UNK variable that is due to
    particles, or leaves "pden" as it is, after a Poisson equation solve.
    This only applies meaningfully when a "pden" variable is declared and
    the gravitational potential is calculated by solving a Poisson equation
    whose right-hand side includes a mass distribution to which both
    hydrodynamic fluid density and massive particles contribute. The "pden"
    variable will have been set to the sum of the fluid density ("dens"
    variable) and the density resulting from mapping massive particles to
    the mesh, so that is what remains in "pden" when grav_unjunkPden is set
    to FALSE. Otherwise, "dens" will be subtracted from "pden" before
    Gravity_potentialListOfBlocks returns, and "pden" will be left
    containing only the mass density that is due to particles.
  point_mass [REAL] [0.e0]
    Valid Values: Unconstrained
    mass of the central point-like object
  point_mass_rsoft [REAL] [0.e0]
    Valid Values: Unconstrained
    softening radius for the point-like mass (in units of number of the
    finest level cells)
  updateGravity [BOOLEAN] [TRUE]
    allow gravity value to be updated
```

# Equation of States

- **Gamma-law equation of state could be invalid in some astrophysical environments in which electrons and positrons may be relativistic and/or degenerate, and in which radiation may significantly contribute to the thermodynamic state.**
- **The Helmholtz EOS includes contributions from radiation, completely ionized nuclei, and degenerate/relativistic electrons and positrons. The pressure and internal energy are calculated as the sum over the components.**

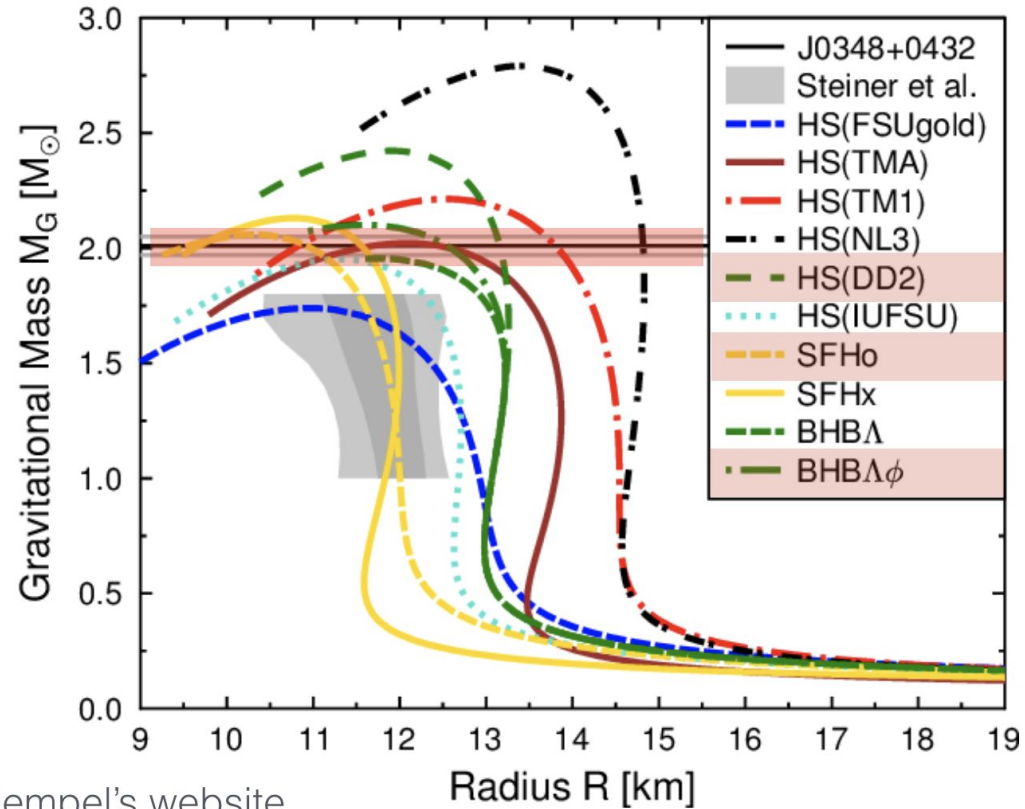
$$P_{\text{tot}} = P_{\text{rad}} + P_{\text{ion}} + P_{\text{ele}} + P_{\text{pos}} + P_{\text{coul}}$$

# Equation of States

- **Extreme hot and dense environments can be realized in core-collapse supernovae, which occur at the very end of massive stellar evolutions and lead to the formation of a neutron star or a black hole.**
- **At high temperature  $T > 0.4 \text{ MeV}$ , chemical equilibrium is achieved for all strong and electromagnetic reactions, which is referred to as nuclear statistical equilibrium (NSE) and the nuclear composition is determined as a function of density, temperature, and proton (electron) fraction.**

$$10^5 < \rho < 10^{15} \text{ g/cm}^3 \quad 0.1 < T < 100 \text{ MeV}$$

# Nuclear EoS



From M. [Hempel's](#) website

# Multi-species

source/Simulation/SimulationComposition/Burn

1	neut	0	1	1	0	2	0.5			
2	h1	1		1		0		0	2	0.5
3	prot	1	1	0		0	2	0.5		
4	he3	2	3	1	7.718	2	0			
5	he4	2	4	2		28.296	1	0		
6	c12	6	12	6		92.162	1	0		
7	n14		7		14		7	104.659	1	0
8	o16	8	16	8		127.619	1	0		
9	ne20		10	20	10		160.645	1	0	
10	ne21		10	21	11		167.406	4	0	
11	ne22		10	22	12		177.77	1	0	
12	na23		11	23	12		186.564	4	0	
13	mg23		12	23	11		181.726	4	0	
14	mg24		12	24	12		198.257	1	0	
15	mg25		12	25	13		205.588	6	0	
16	mg26		12	26	14		216.681	1	0	
17	al25		13	25	12		200.529	6	0	
18	al26		13	26	13		211.894	11	0	

< Example: species in Burn >

# Modify your simulation

Simulation: `Simulation_adjustSimulation(blkcnt, blklst, nstep, dt, stime)`

Physics

Heating: `Heat(blockCount,blockList,dt,time)`

Cooling: `Cool(blockCount,blockList,dt,time)`

**< Example: nuclear heating in Project A >**

# Modify AMR

API: Grid\_markRefineDerefine(), Grid\_markRefineSpecialized

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## SYNOPSIS

```
Grid_markRefineSpecialized(integer(IN) :: criterion,  
                           integer(IN) :: size,  
                           real(IN)    :: specs(size),  
                           integer(IN) :: lref )
```

# Modify AMR

## DESCRIPTION

The routine provides an interface to a collection of routines that define very specialized refinement criteria. The currently supported options are:

- THRESHOLD : when a specific variable is below or above a threshold
- ELLIPSOID : The blocks that fall within the specified ellipsoid
- RECTANGLE : The blocks that fall within the specified rectangle
- INRADIUS : The blocks that fall within the specified radius
- WITHRADIUS : The blocks that fall on the specified radius



# Modify AMR

## ARGUMENTS

```

criterion - the creterion on which to refine
size      - size of the specs data structure
specs     - the data structure containing information specific to
            the creterion
For THRESHOLD
    specs(1) = real(variable_name), for example
              if variable is density, then
              specs(1)=real(DENS_VAR)
    specs(2) = the threshold value
    specs(3) = if < 0 refine if variable < threshold
              if > 0 refine if variable > threshold
For ELLIPSOID
    specs(1:3) = center of the ellipsoid
    specs(4:6) = the semimajor axes of the ellipsoid
For INRADIUS
    specs(1:3) = center of the circle/sphere
    specs(4)   = the radius
For WITHRADIUS
    specs(1:3) = center of the circle/sphere
    specs(4)   = the radius
For RECTANGLE
    specs(1:6) = bounding coordinates of rectangle
    specs(7)   = if 0 refine block with any overlap
              if /= refine only blocks fully
              contained in the rectangle

lref      - If > 0, bring selected blocks to this level of refinement.
            If <= 0, refine qualifying blocks once.
```